
Speed of Sound In QCD Plasma

Rajiv Gavai, Sourendu Gupta, Swagato Mukherjee

Tata Institute of Fundamental Research
Mumbai

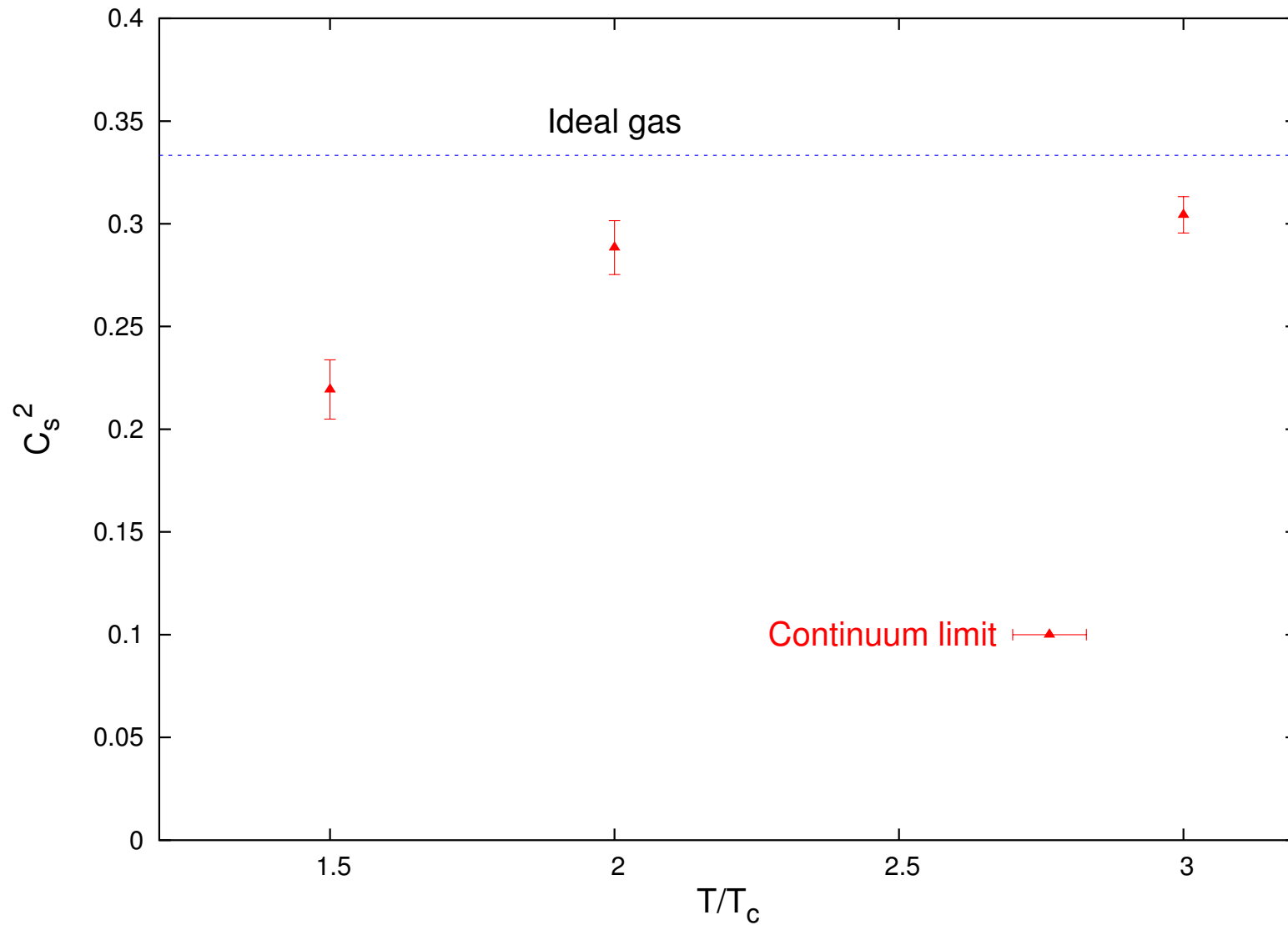
Introduction

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- The QCD plasma, created in heavy-ion collisions, may contain long wavelength oscillations, like sound waves, near equilibrium if its viscosity is small enough.
- These sound waves in a QCD plasma near equilibrium may affect its particle spectrum and can be observed.

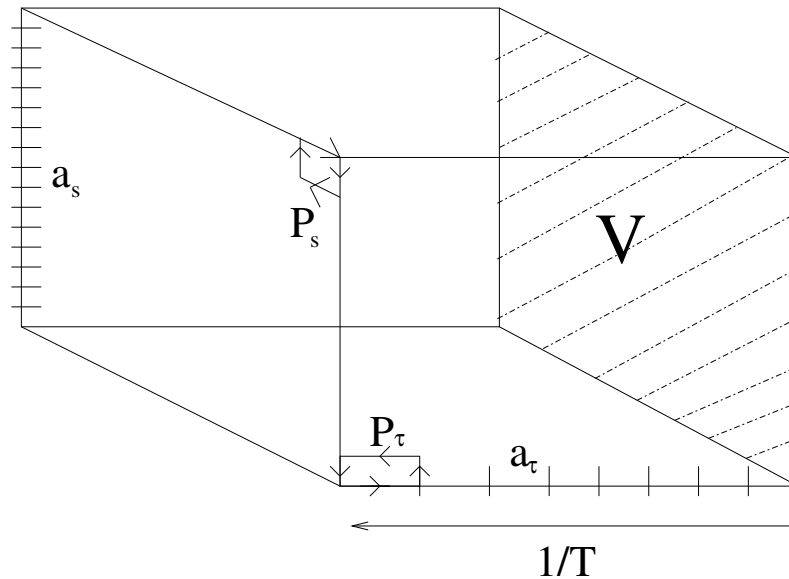
Speed of sound



Continuum limit of speed of sound in plasma phase of quenched QCD theory.

Action

- Wilson action on an asymmetric lattice :



$$S(U) = 2N_c K_s P_s + 2N_c K_\tau P_\tau$$

$$K_s \equiv \frac{1}{\xi g_s^2} \quad K_\tau \equiv \frac{\xi}{g_\tau^2} \quad \xi = \frac{a_s}{a_\tau}$$

Thermodynamics

- Partition function :

$$\mathcal{Z} = \int \mathcal{D}U e^{-S(U)}$$

- Energy density :

$$E = \frac{T^2}{V} \frac{\partial \ln \mathcal{Z}}{\partial T} \Big|_V$$

- Pressure :

$$P = T \frac{\partial \ln \mathcal{Z}}{\partial V} \Big|_T$$

Energy density

- Energy density :

$$\frac{E}{T^4} = 6N_c N_\tau^4 \left[\frac{\Delta_s - \Delta_\tau}{g^2} - (c'_s \Delta_s + c'_\tau \Delta_\tau) \right]$$

where :

$$\Delta_j = \langle \bar{P}_j \rangle - \langle \bar{P}_0 \rangle$$

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- Karsch coefficients, evaluated upto one loop order :

$$g_i^{-2}(a_s, \xi) = g^{-2}(a) + c_i(\xi) + O[g^2(a)]$$

Specific heat

- Specific heat at constant volume :

$$\frac{C_v}{T^3} = \frac{1}{T^3} \frac{\partial E}{\partial T} =$$

$$\frac{4E}{T^4} - 6N_c N_\tau^4 \left[2g^{-2} \Delta_\tau + 4c'_\tau \Delta_\tau + (c''_s \Delta_s + c''_\tau \Delta_\tau) \right] +$$

$$36N_c^2 N_s^3 N_\tau^5 \left[g^{-4} \text{var}(\Delta_s - \Delta_\tau) + \right.$$

$$\left. g^{-2} \text{var}(c'_s \Delta_s + c'_\tau \Delta_\tau, \Delta_s - \Delta_\tau) + \text{var}(c'_s \Delta_s + c'_\tau \Delta_\tau) \right]$$

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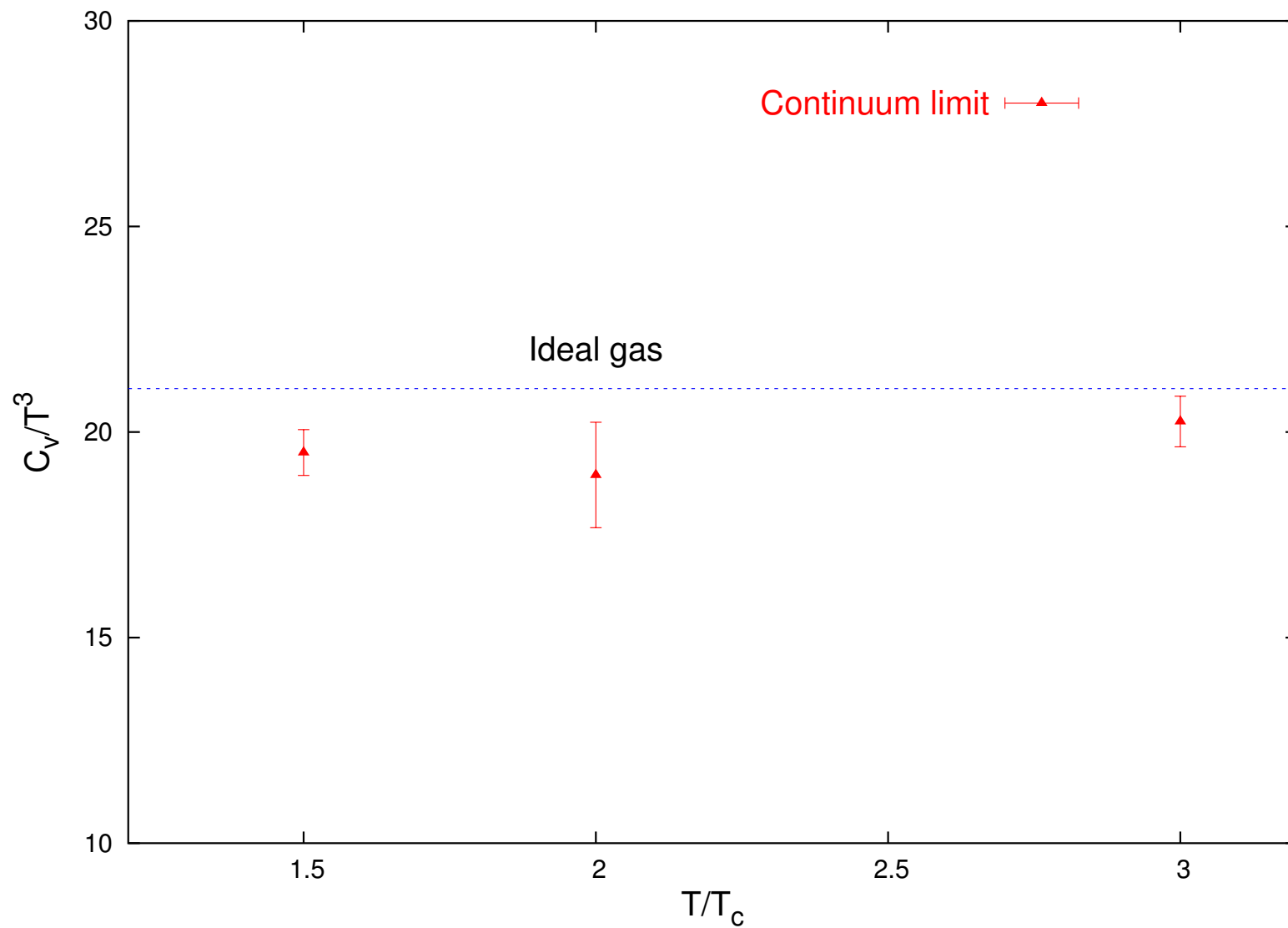
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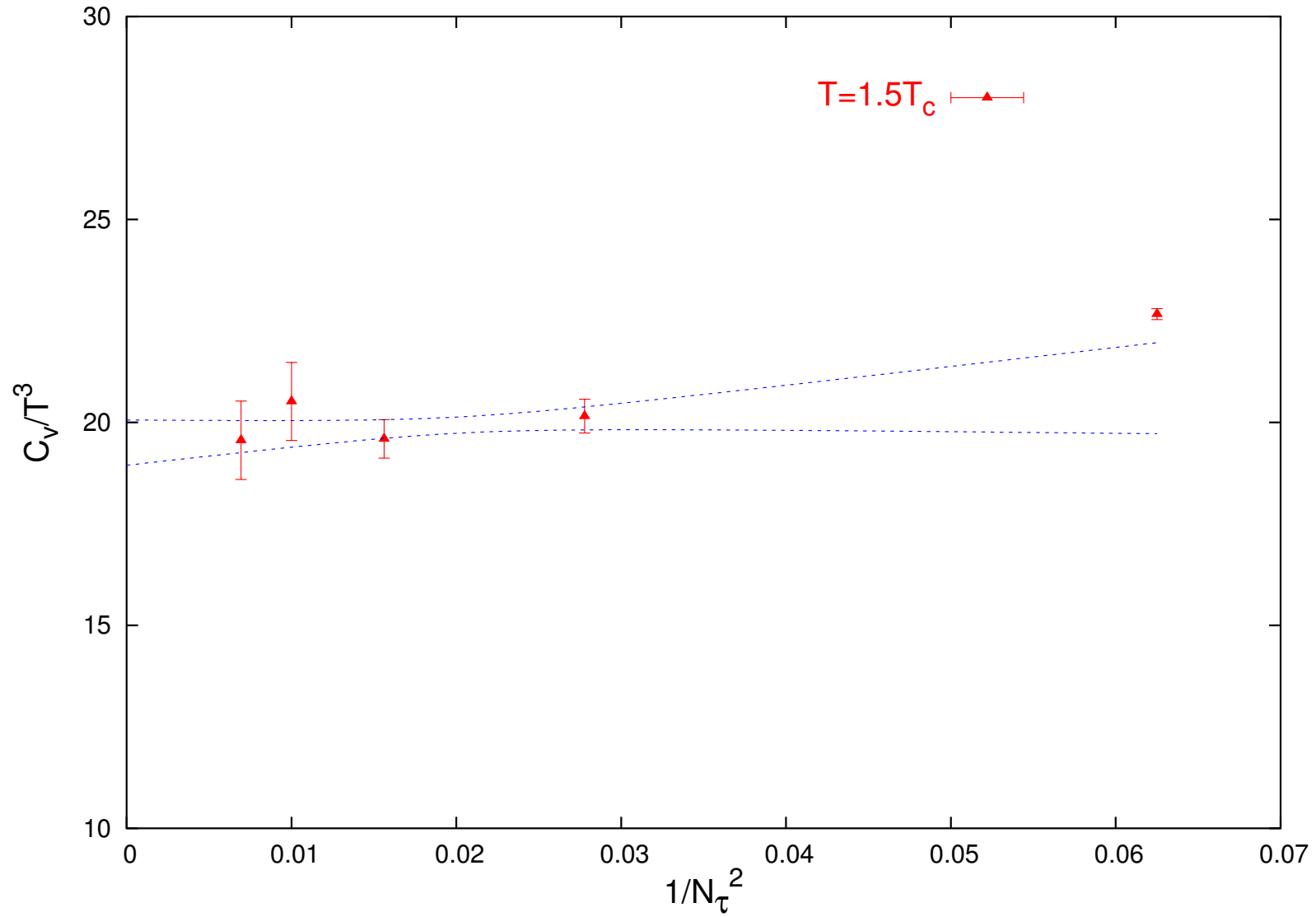
$$\frac{4E}{T^4} - 6N_c N_\tau^4 \left[2g^{-2} \Delta_\tau + 4c'_\tau \Delta_\tau + (c''_s \Delta_s + c''_\tau \Delta_\tau) \right]$$

- $c''_s(\xi = 1) = -0.298192$
 $c''_\tau(\xi = 1) = 0.333674$

C_v Results



C_v Results



Speed of sound

- Speed of sound :

$$C_s^2 = \left. \frac{\partial P}{\partial E} \right|_s = \frac{1}{3} - \frac{1}{3} \cdot \frac{\frac{1}{T^3} \left(\frac{\partial D}{\partial T} \right)_V}{\frac{1}{T^3} \left(\frac{\partial E}{\partial T} \right)_V}$$

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- Interaction measure :

$$\frac{D}{T^4} = \frac{E - 3P}{T^4} = 6N_c N_\tau^4 \left(a \frac{\partial g^{-2}}{\partial a} \right) (\Delta_s + \Delta_\tau)$$

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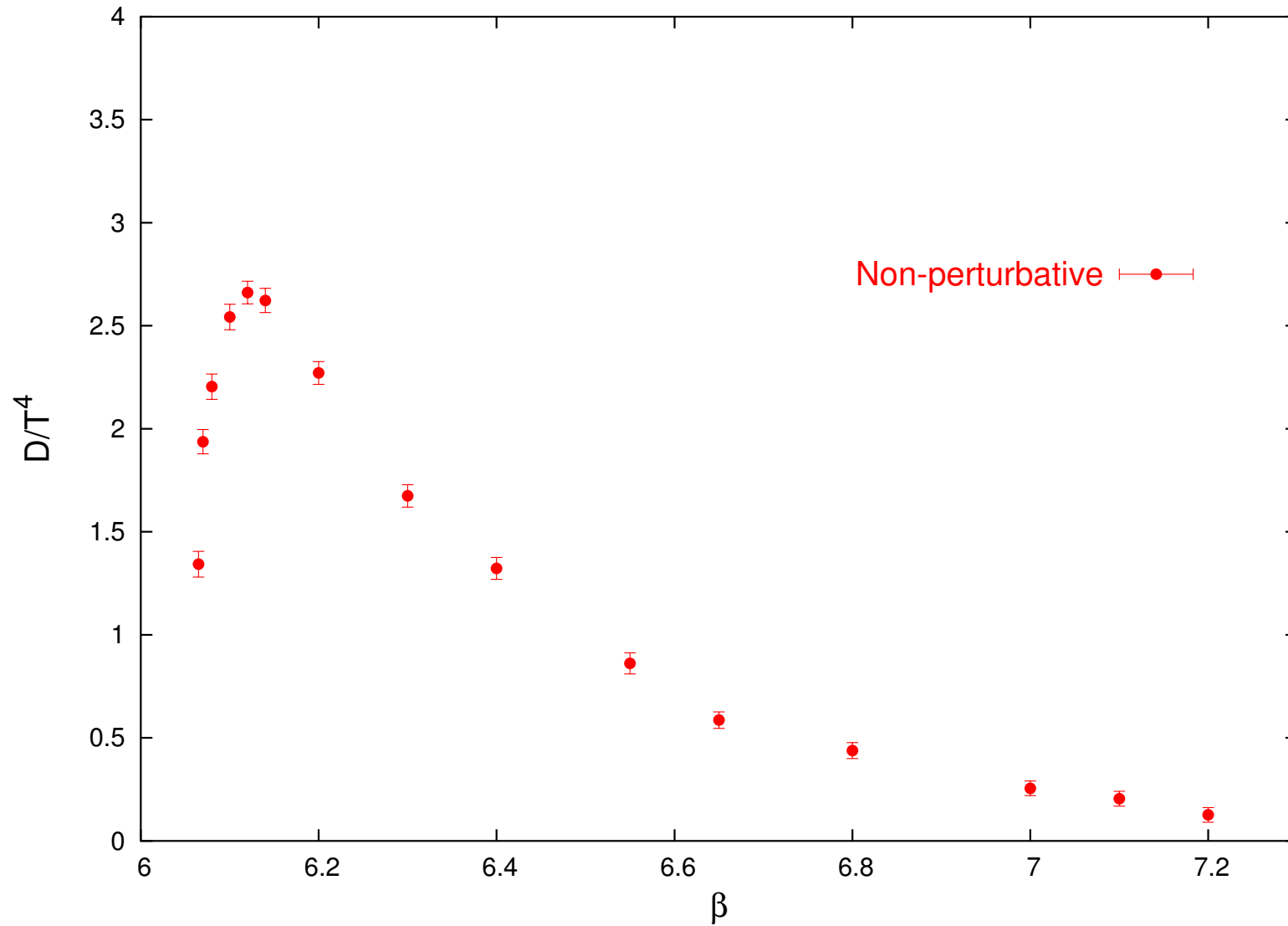
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- Derivative of interaction measure :

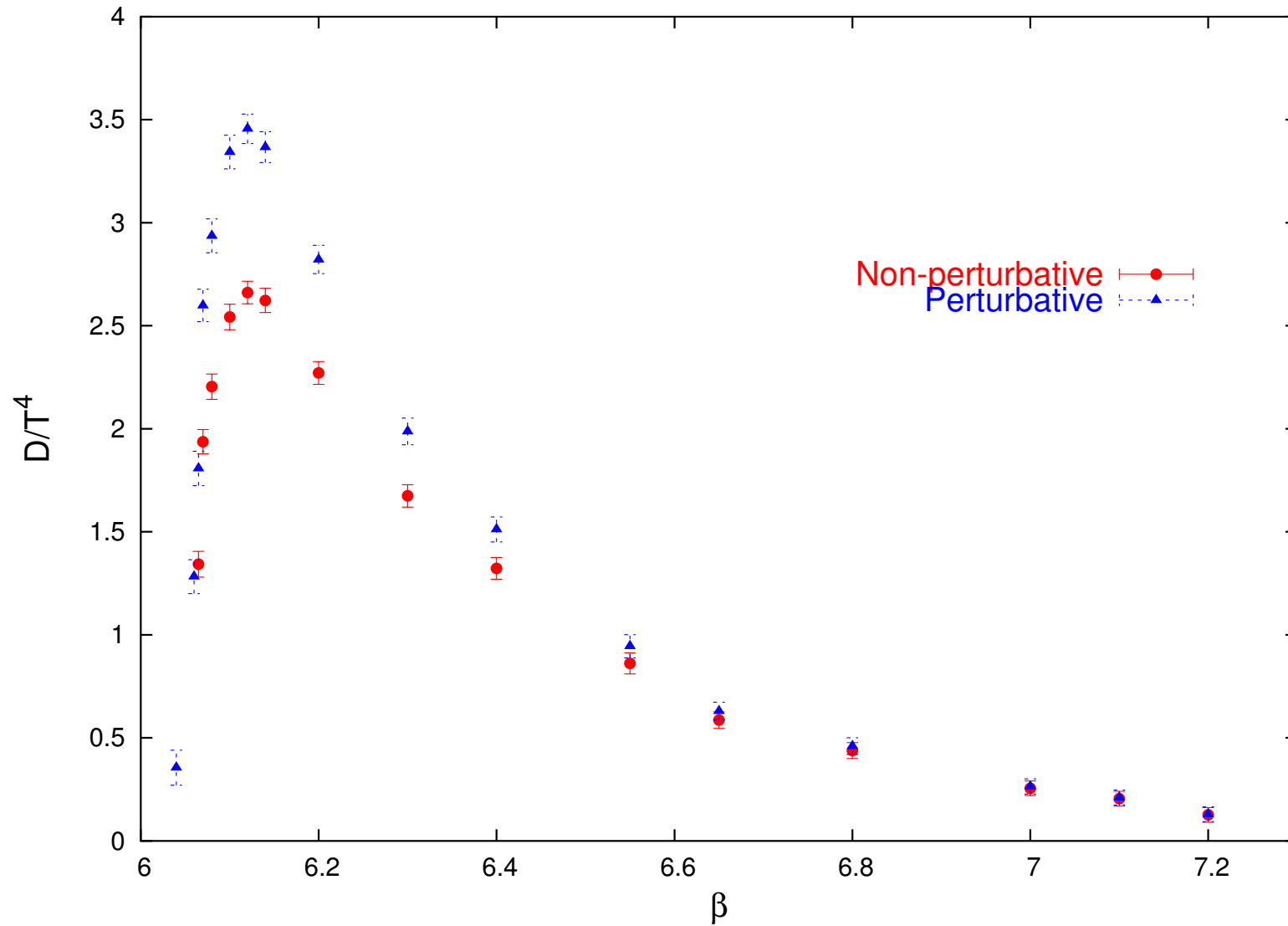
$$\frac{1}{T^3} \frac{\partial D}{\partial T} = \frac{4D}{T^4} - 12N_c N_\tau^4 \left(a \frac{\partial g^{-2}}{\partial a} \right) \Delta_\tau$$

Lattice sizes



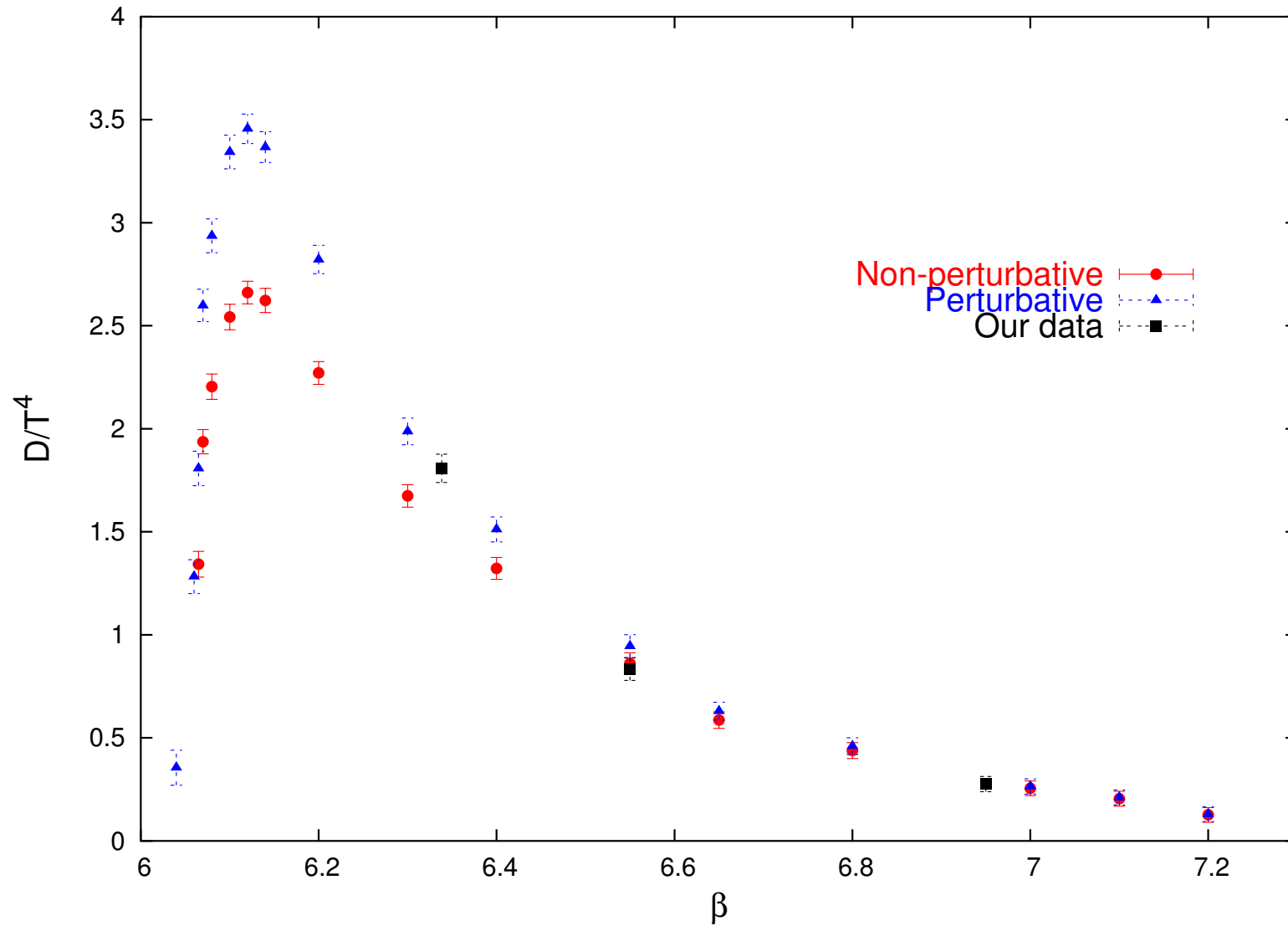
G.Boyd *et al.*, Nucl. Phys., B 469 (1996) 419.

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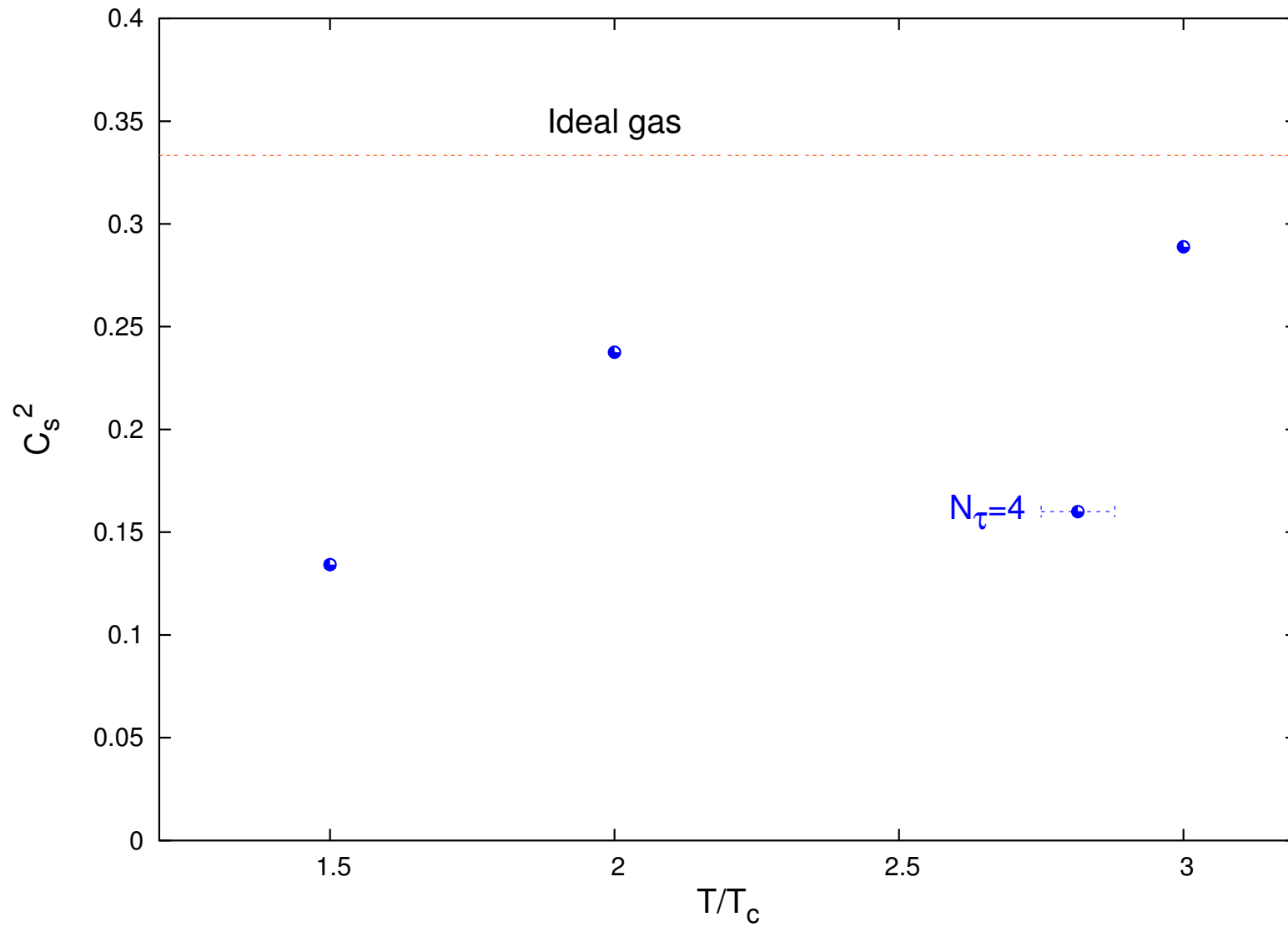
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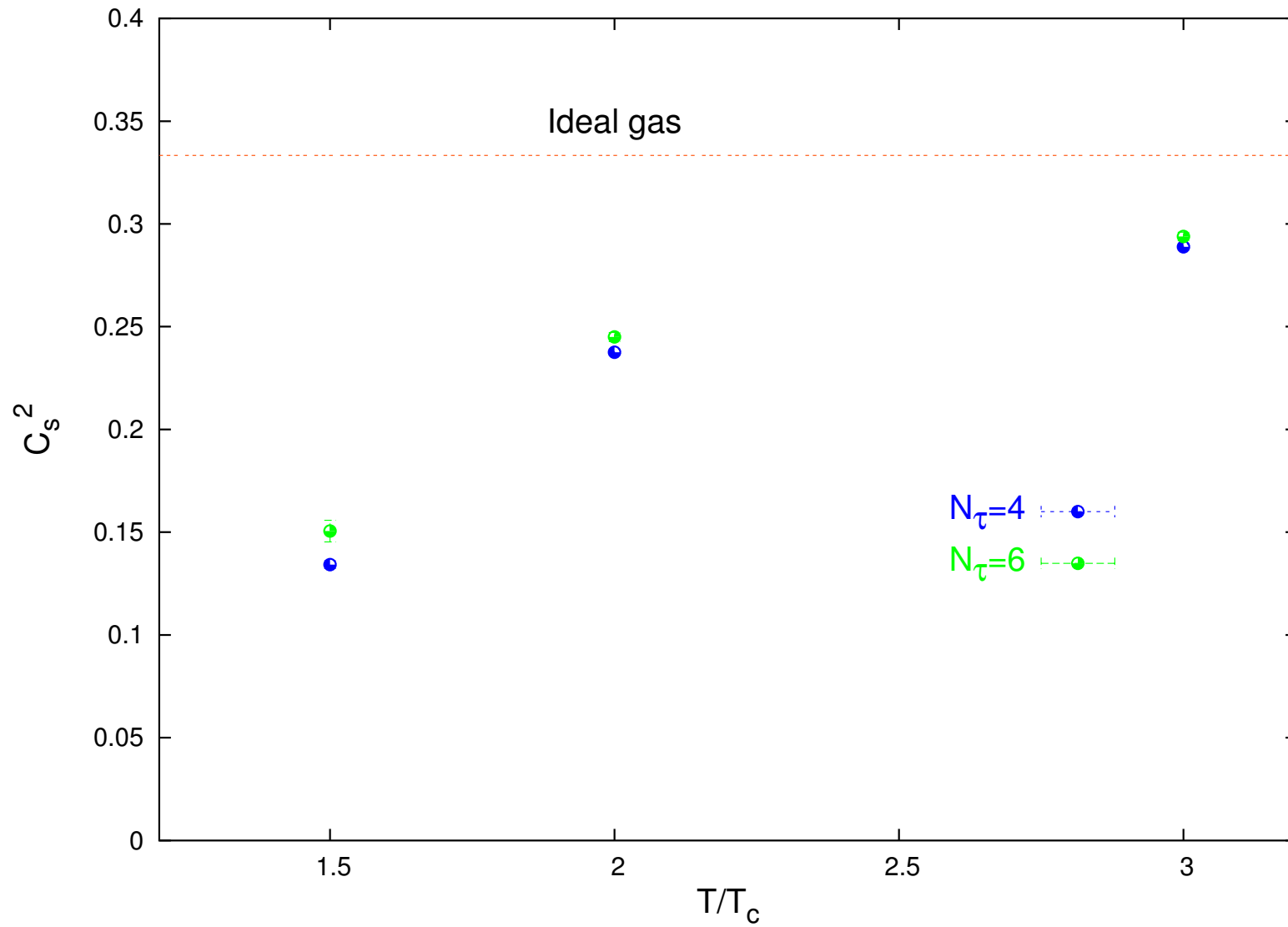


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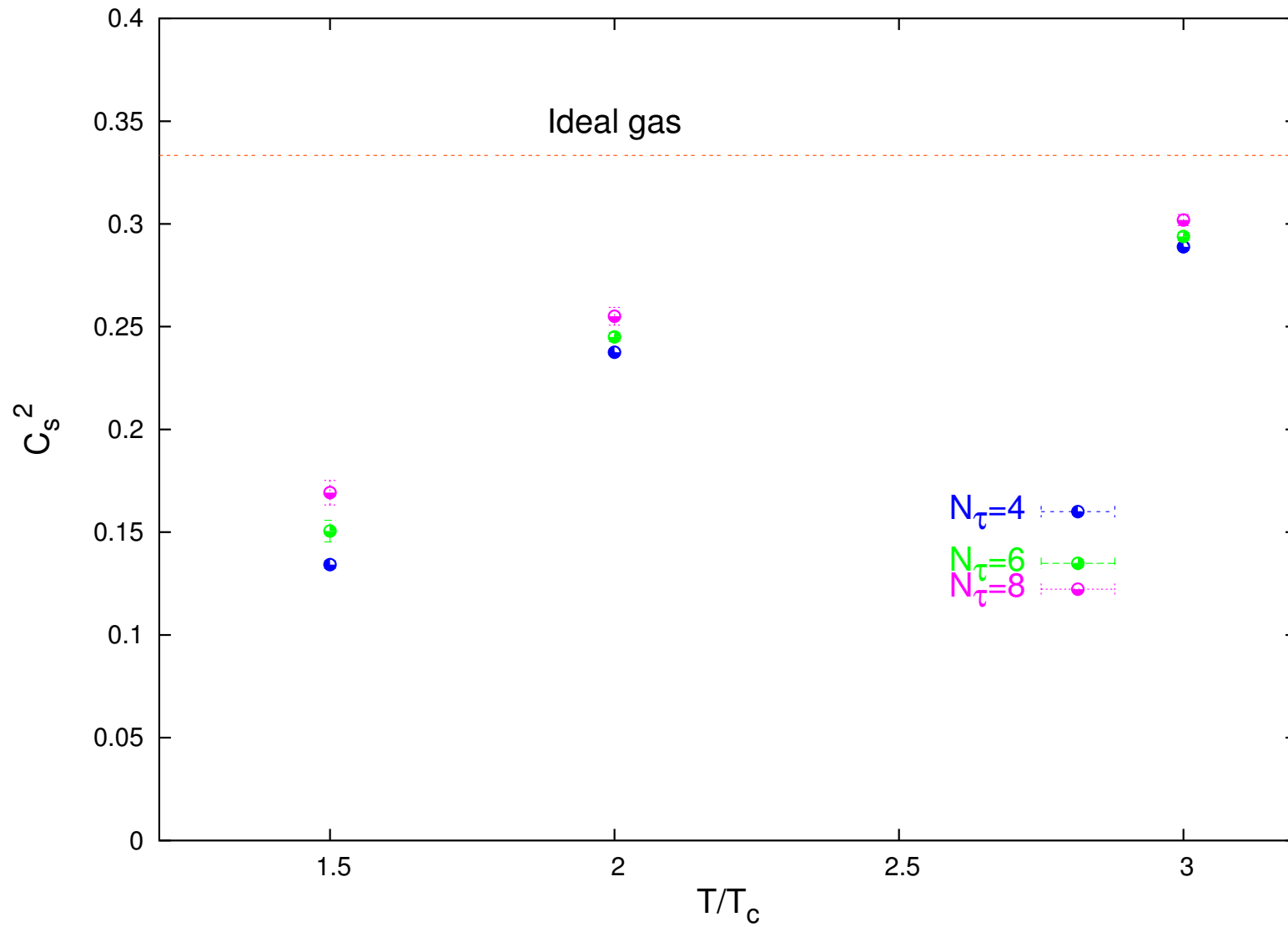
C_s^2 Results



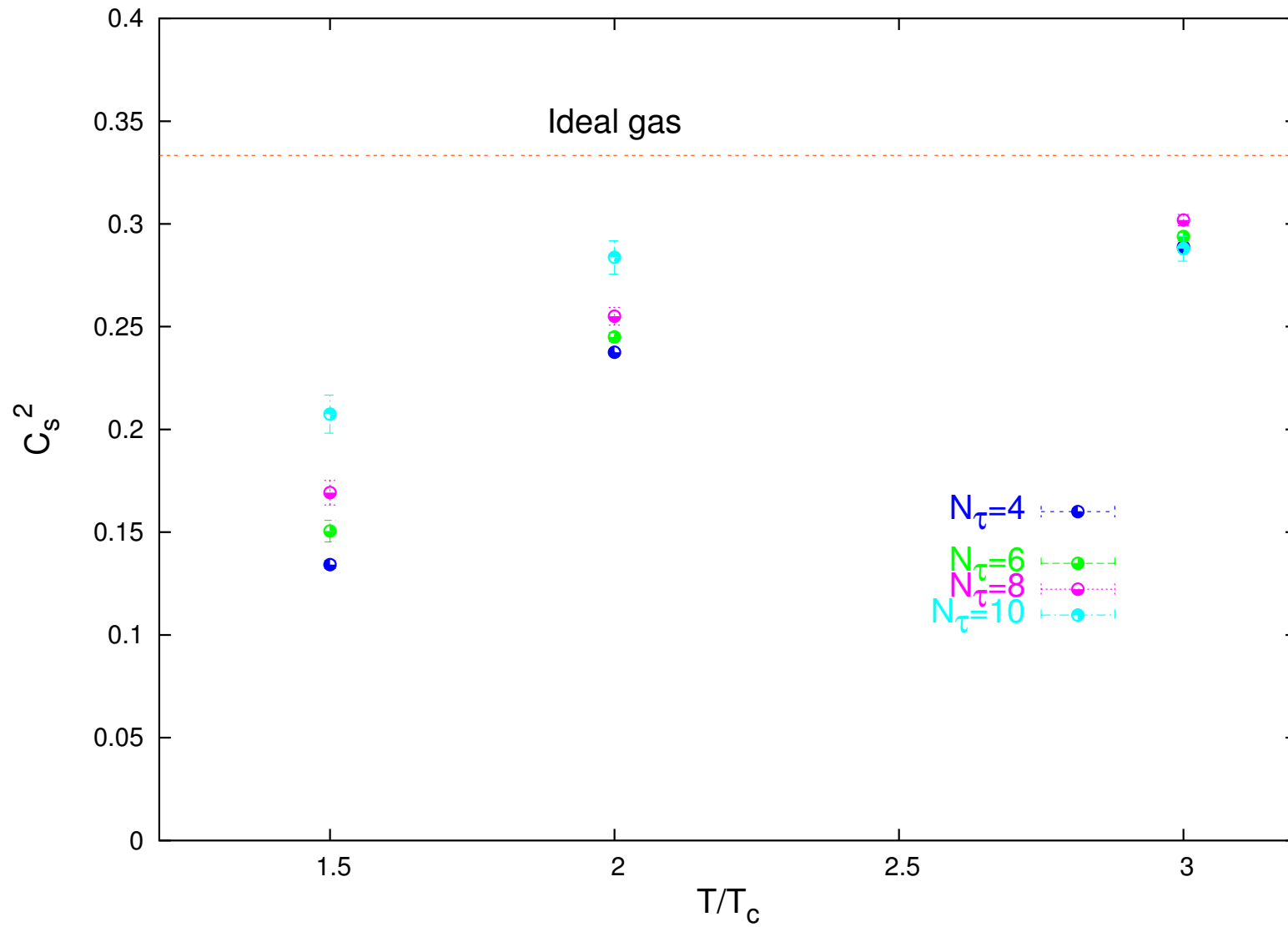
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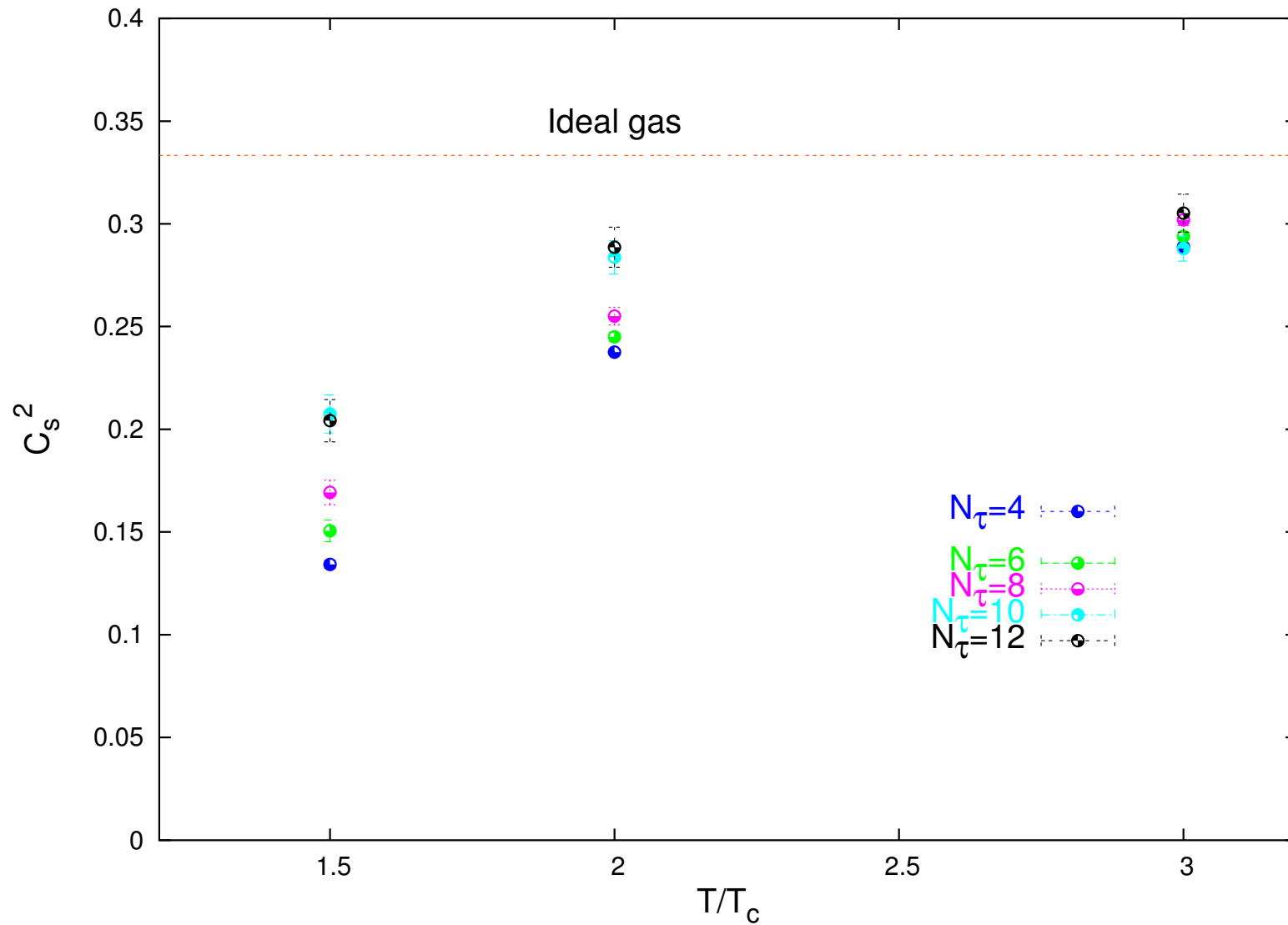
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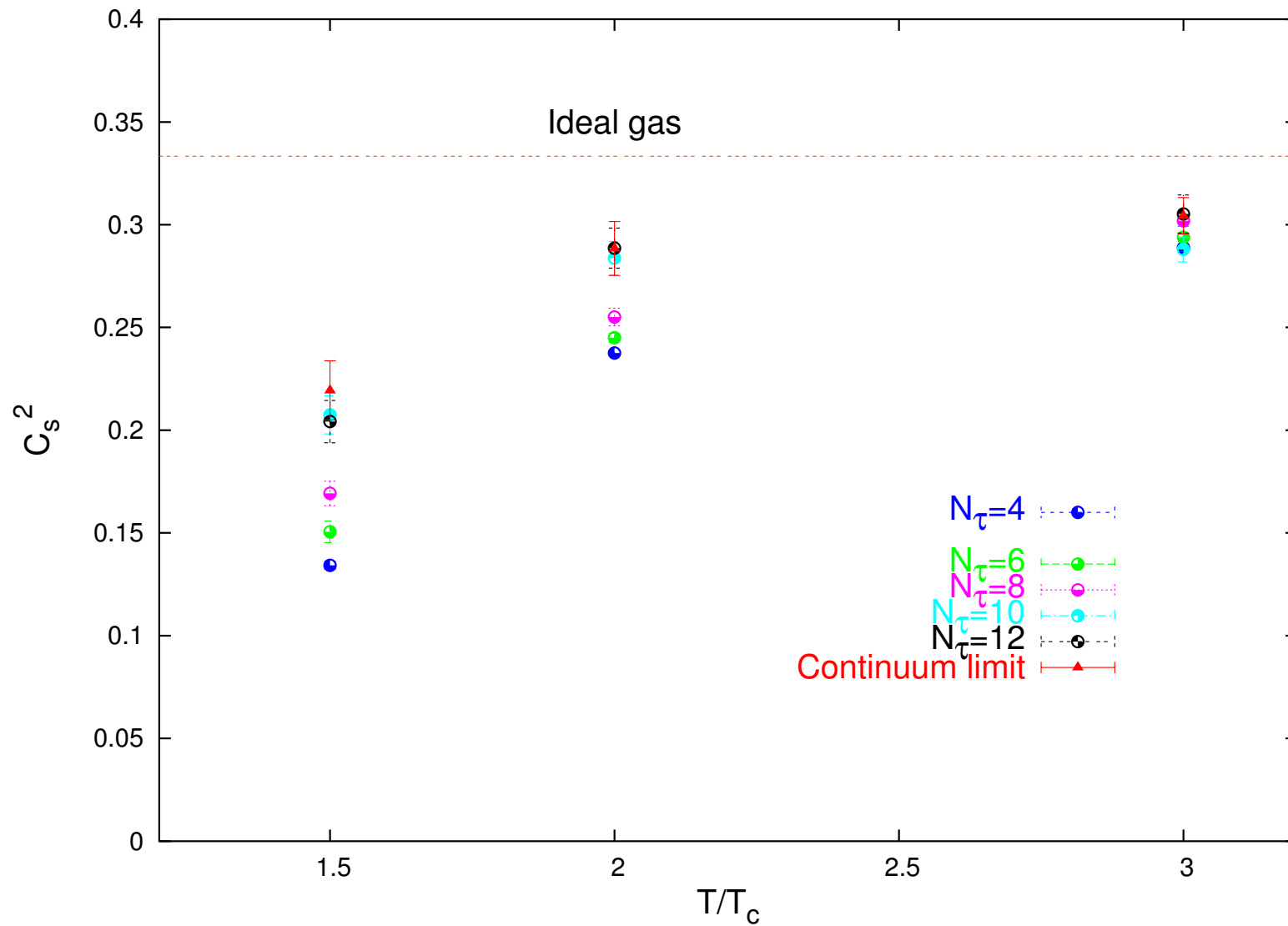
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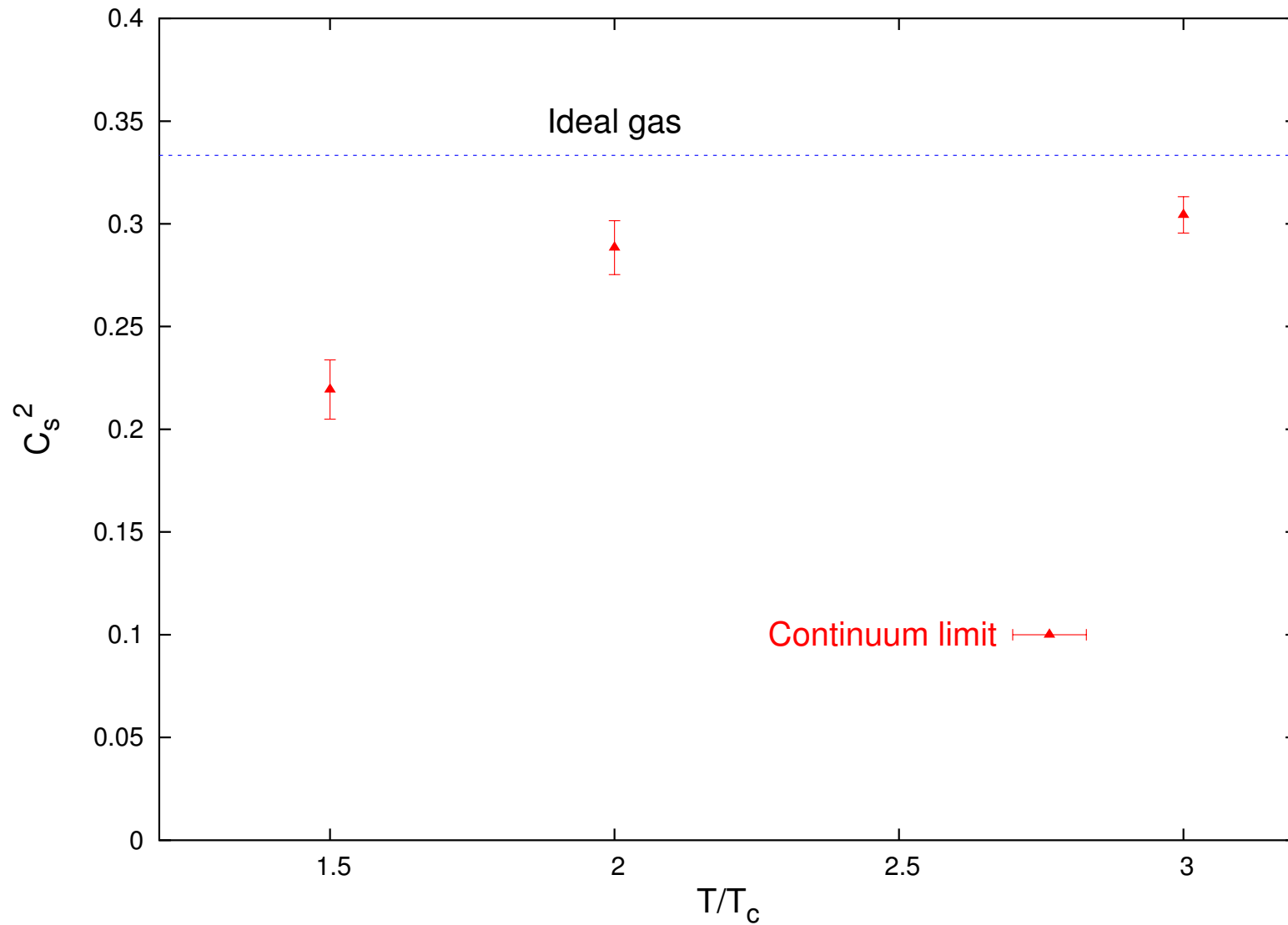
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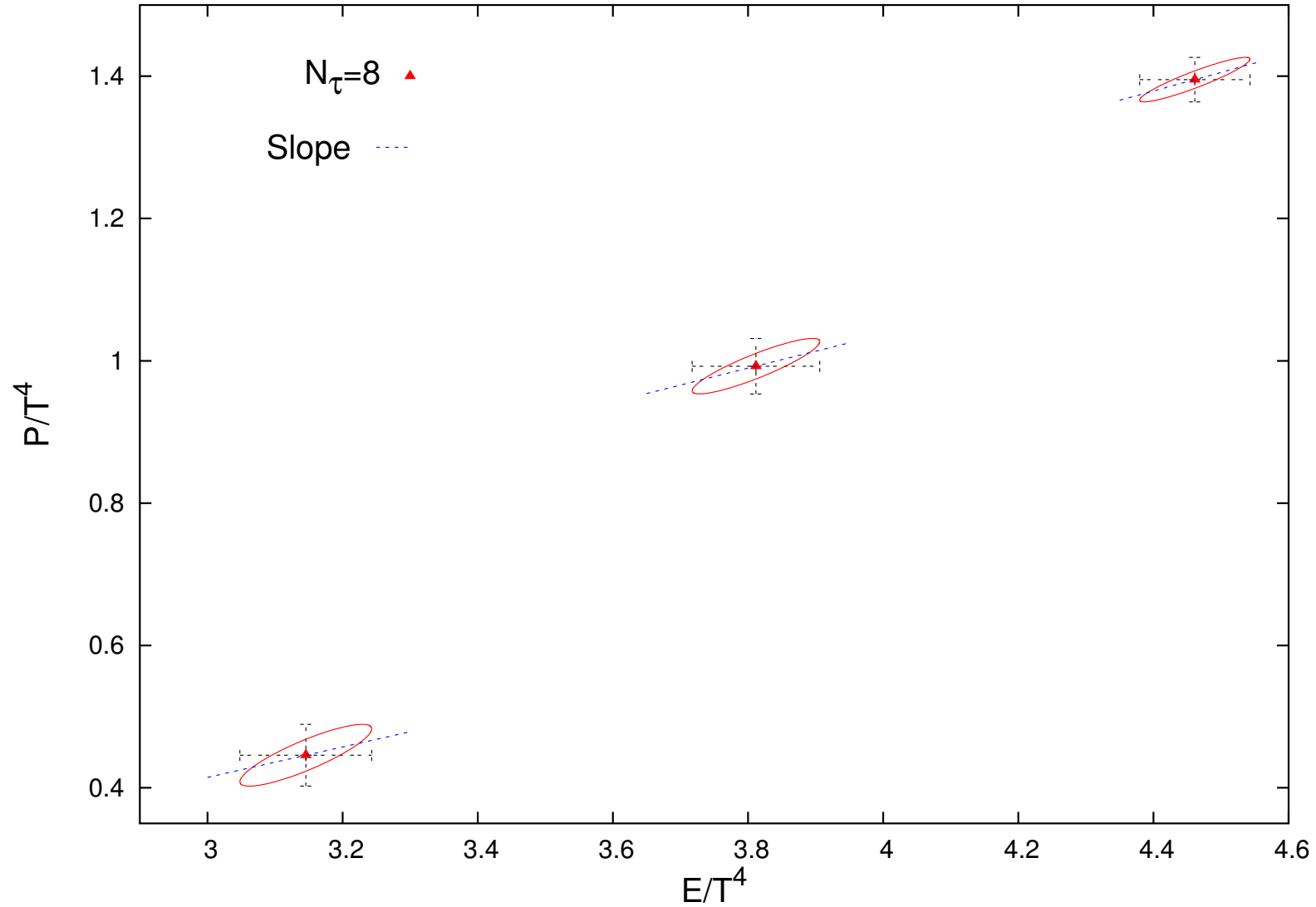
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P vs E



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- At $T = 3T_c$ the continuum limit of C_v differs from its value for the ideal gas by 4% with 80% confidence limit.
- At $T = 3T_c$ the continuum limit of C_s^2 differs from its value for the ideal gas by 9% with 99% confidence limit.